

# A Phenomenology of the Baryon Spectrum from Lattice QCD

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Jefferson Lab

*N<sup>\*</sup> 2011*

Collaborators:

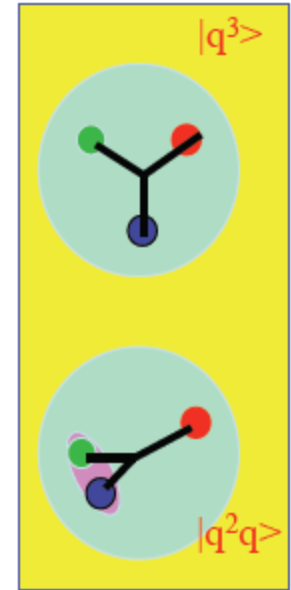
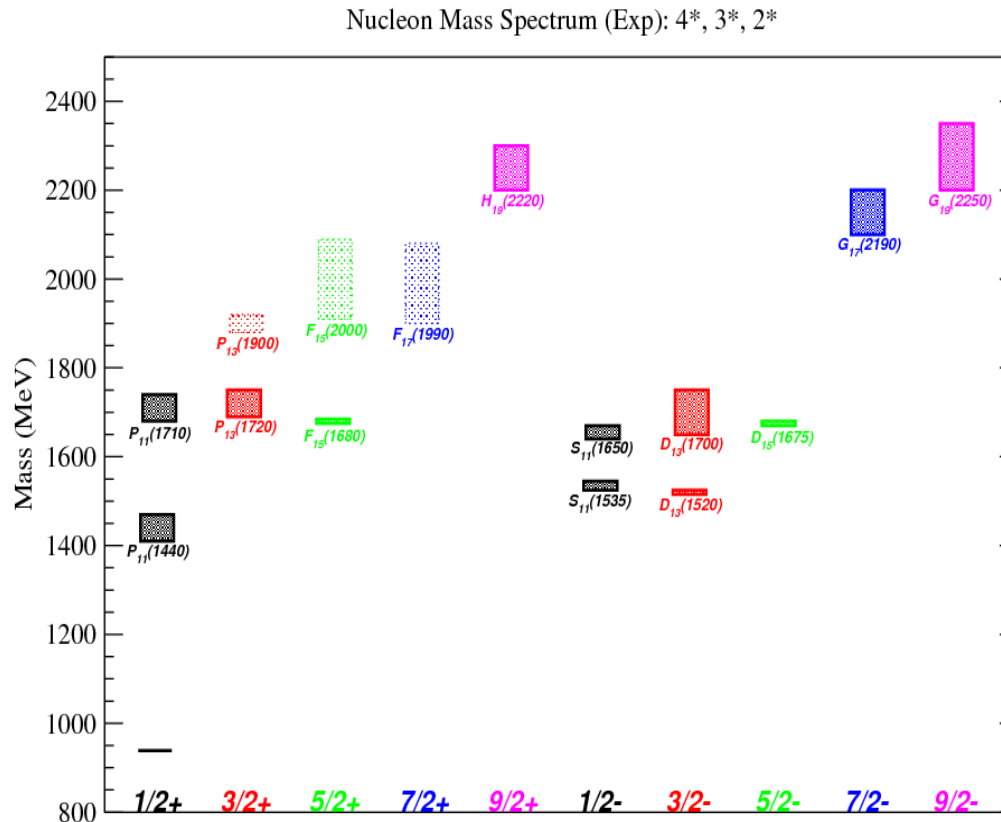
J. Dudek, B. Joo, D. Richards, S. Wallace

Auspices of the Hadron Spectrum Collaboration

# Baryon Spectrum

“Missing resonance problem”

- What are collective modes?
- What is the structure of the states?



Nucleon spectrum

PDG uncertainty on B-W mass

# Lattice QCD

**Goal:** resolve highly excited states

$$N_f = 2 + 1 \text{ (u,d + s)}$$

**Anisotropic lattices:**

$$(a_s)^{-1} \sim 1.6 \text{ GeV}, \quad (a_t)^{-1} \sim 5.6 \text{ GeV}$$

0810.3588, 0909.0200, 1004.4930

# Spectrum from variational method

Two-point correlator

$$C(t) = \langle 0 | \Phi'(t) \Phi(0) | 0 \rangle$$

$$C(t) = \sum_{\mathbf{n}} e^{-E_{\mathbf{n}} t} \langle 0 | \Phi'(0) | \mathbf{n} \rangle \langle \mathbf{n} | \Phi(0) | 0 \rangle$$

$$Z_i^n \equiv \langle \mathbf{n} | \Phi_i | 0 \rangle$$

Matrix of correlators

$$C(t) = \begin{bmatrix} \langle 0 | \Phi_1(t) \Phi_1(0) | 0 \rangle & \langle 0 | \Phi_1(t) \Phi_2(0) | 0 \rangle & \dots \\ \langle 0 | \Phi_2(t) \Phi_1(0) | 0 \rangle & \langle 0 | \Phi_2(t) \Phi_2(0) | 0 \rangle & \dots \\ \vdots & & \ddots \end{bmatrix}$$

Diagonalize:

eigenvalues  $\rightarrow$  spectrum

eigenvectors  $\rightarrow$  spectral “overlaps”

Each state optimal combination of  $\Phi_i$

$$\Omega_{\mathbf{n}} = v_1^n \Phi_1 + v_2^n \Phi_2 + \dots$$

Benefit: orthogonality for near degenerate states

# Operator construction

Baryons : permutations of 3 objects

Permutation group  $S_3$ : 3 representations

- **Symmetric**: 1-dimensional
  - e.g.,  $uud+udu+duu$
- **Antisymmetric**: 1-dimensional
  - e.g.,  $uud-udu+duu-\dots$
- **Mixed**: 2-dimensional
  - e.g.,  $udu - duu$  &  $2duu - udu - uud$

Color antisymmetric  $\rightarrow$  Require **Space** [**Flavor Spin**] symmetric

Classify operators by these permutation symmetries:

- Leads to rich structure

1104.5152

# Orbital angular momentum via derivatives

Couple derivatives onto single-site spinors:  
Enough D's – build any J,M

$$\mathcal{O}^{JM} \leftarrow (CGC's)_{i,j,k} \left[ \vec{D} \right]_i \left[ \vec{D} \right]_j [\Psi]_k$$

Only using **symmetries** of continuum QCD

Operator<sub>S</sub>  $\leftarrow$  Derivatives  $\left[ \begin{array}{cc} \text{Flavor} & \text{Dirac} \end{array} \right]$

Use all possible **operators** up to 2 derivatives  
(transforms like 2 units orbital angular momentum)

1104.5152

# Baryon operator basis

3-quark operators with up to two covariant derivatives –  
projected into definite isospin and continuum  $J^P$

$$\text{Operator}_S \leftarrow \left( \left[ \text{Flavor} \quad \text{Dirac} \right] \text{Space}_{\text{symmetry}} \right)^{J^P}$$

Spatial symmetry classification:

$$\text{Nucleons: } N^{2S+1}L_{\pi} J^P$$

Symmetry crucial for spectroscopy

By far the largest operator basis ever used for  
such calculations

$J^P$	#ops	E.g., spatial symmetries	
$J=1/2^-$	24	$N^2P_M \frac{1}{2}^-$	$N^4P_M \frac{1}{2}^-$
$J=3/2^-$	28	$N^2P_M 3/2^-$	$N^4P_M 3/2^-$
$J=5/2^-$	16	$N^4P_M 5/2^-$	
$J=1/2^+$	24	$N^2S_S \frac{1}{2}^+$ $N^2S_M \frac{1}{2}^+$	$N^4D_M \frac{1}{2}^+$ $N^2P_A \frac{1}{2}^+$
$J=3/2^+$	28	$N^2D_S 3/2^+$ $N^2D_M 3/2^+$ $N^2P_A 3/2^+$	$N^4S_M 3/2^+$ $N^4D_M 3/2^+$
$J=5/2^+$	16	$N^2D_S 5/2^+$ $N^2D_M 5/2^+$	$N^4D_M 5/2^+$
$J=7/2^+$	4	$N^4D_M 7/2^+$	

# Operators are not states

Two-point correlator

$$C(t) = \langle 0 | \Phi'(t) \Phi(0) | 0 \rangle$$

$$C(t) = \sum_{\mathbf{n}} e^{-E_{\mathbf{n}} t} \langle 0 | \Phi'(0) | \mathbf{n} \rangle \langle \mathbf{n} | \Phi(0) | 0 \rangle$$

Full basis of operators: many operators can create same state

Spectral “overlaps”

$$\langle \mathbf{n}; J^P | \Phi_i | 0 \rangle = Z_i^{\mathbf{n}}$$

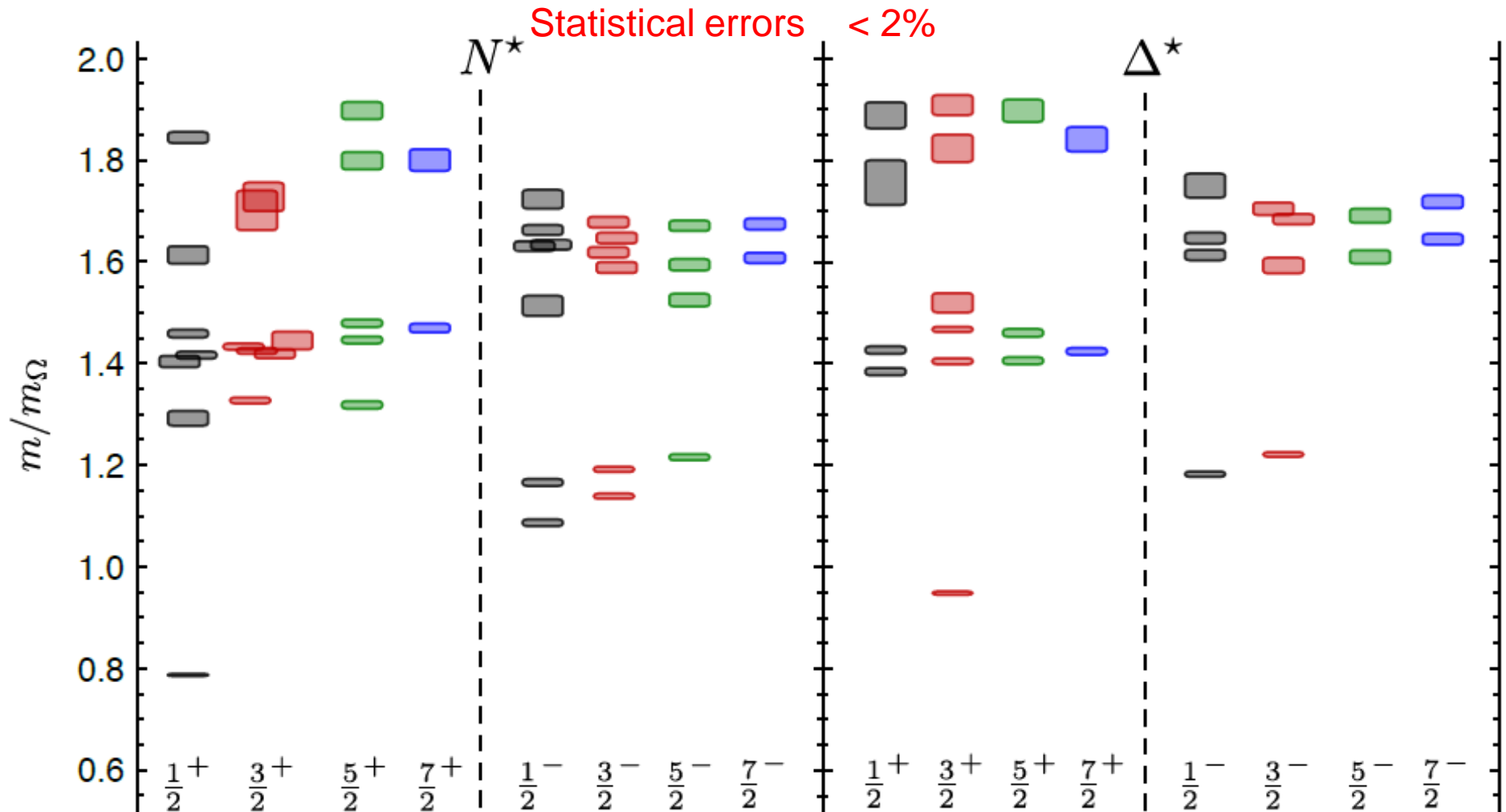
States may have subset of allowed symmetries



# Spin identified Nucleon & Delta spectrum

arXiv:1104.5152

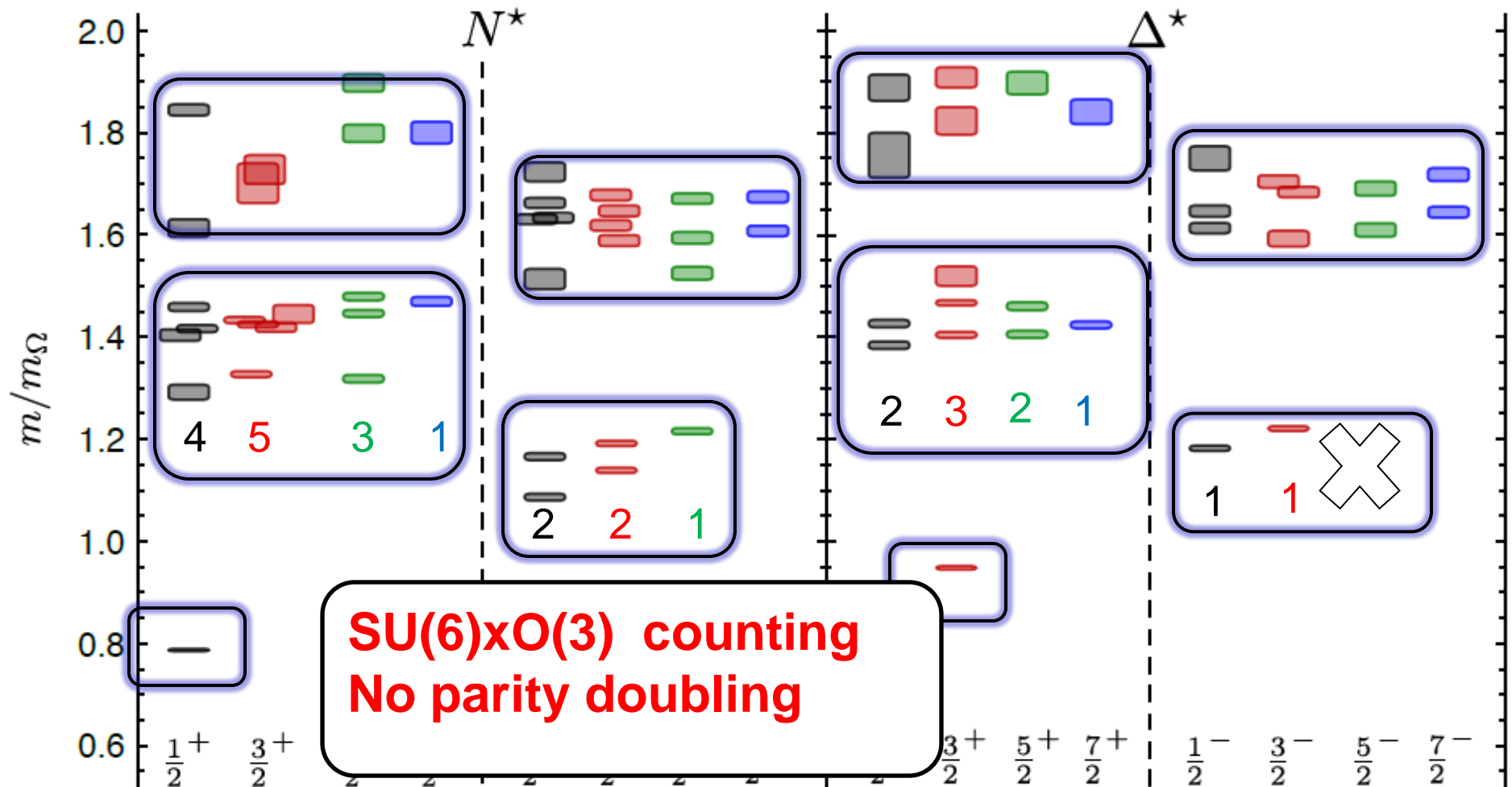
$m_\pi \sim 520\text{MeV}$



# Spin identified Nucleon & Delta spectrum

arXiv:1104.5152

$m_\pi \sim 520\text{MeV}$

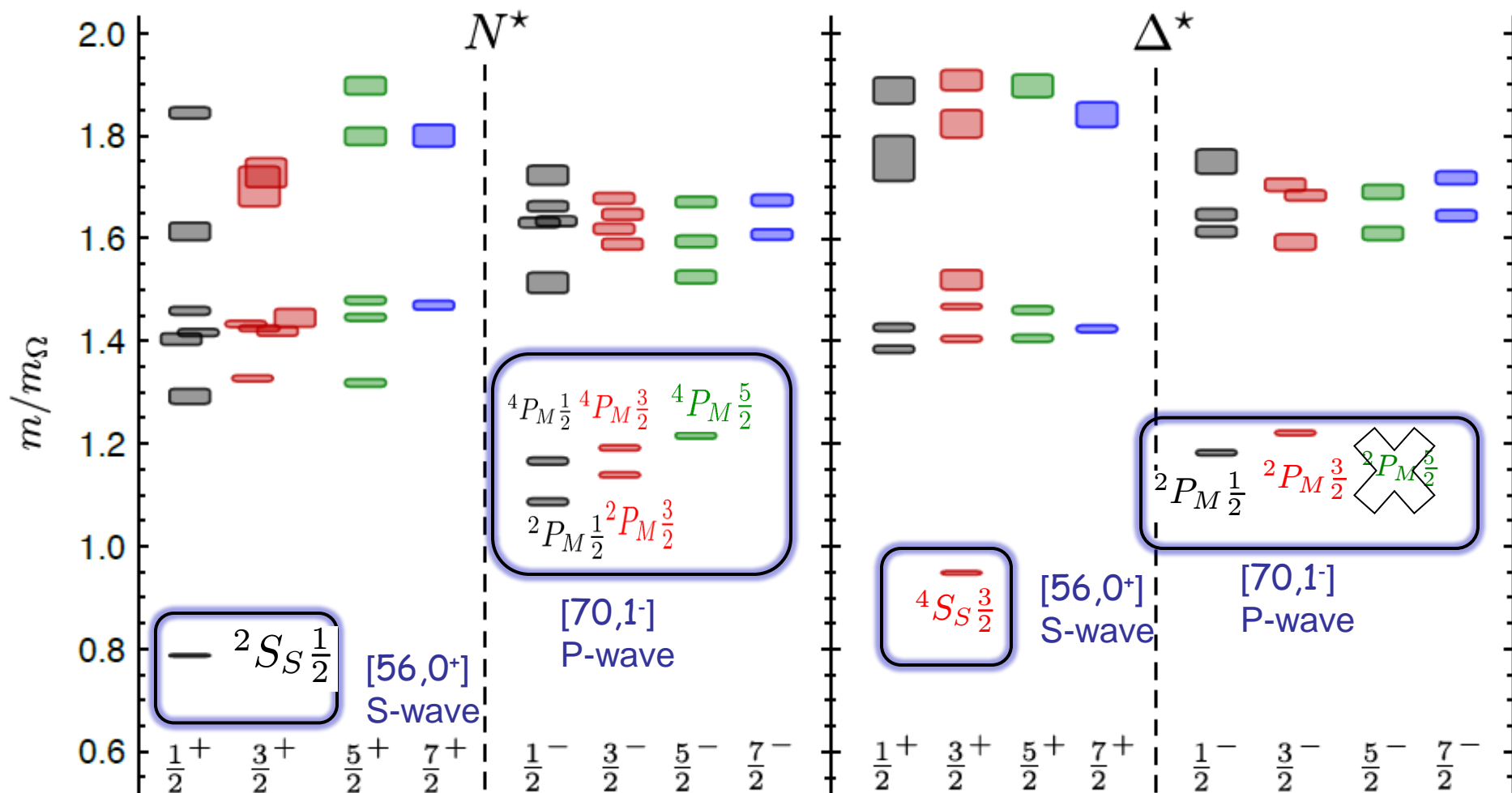


# Spin identified Nucleon & Delta spectrum

Discern structure: spectral overlaps

arXiv:1104.5152

$m_\pi \sim 520\text{MeV}$

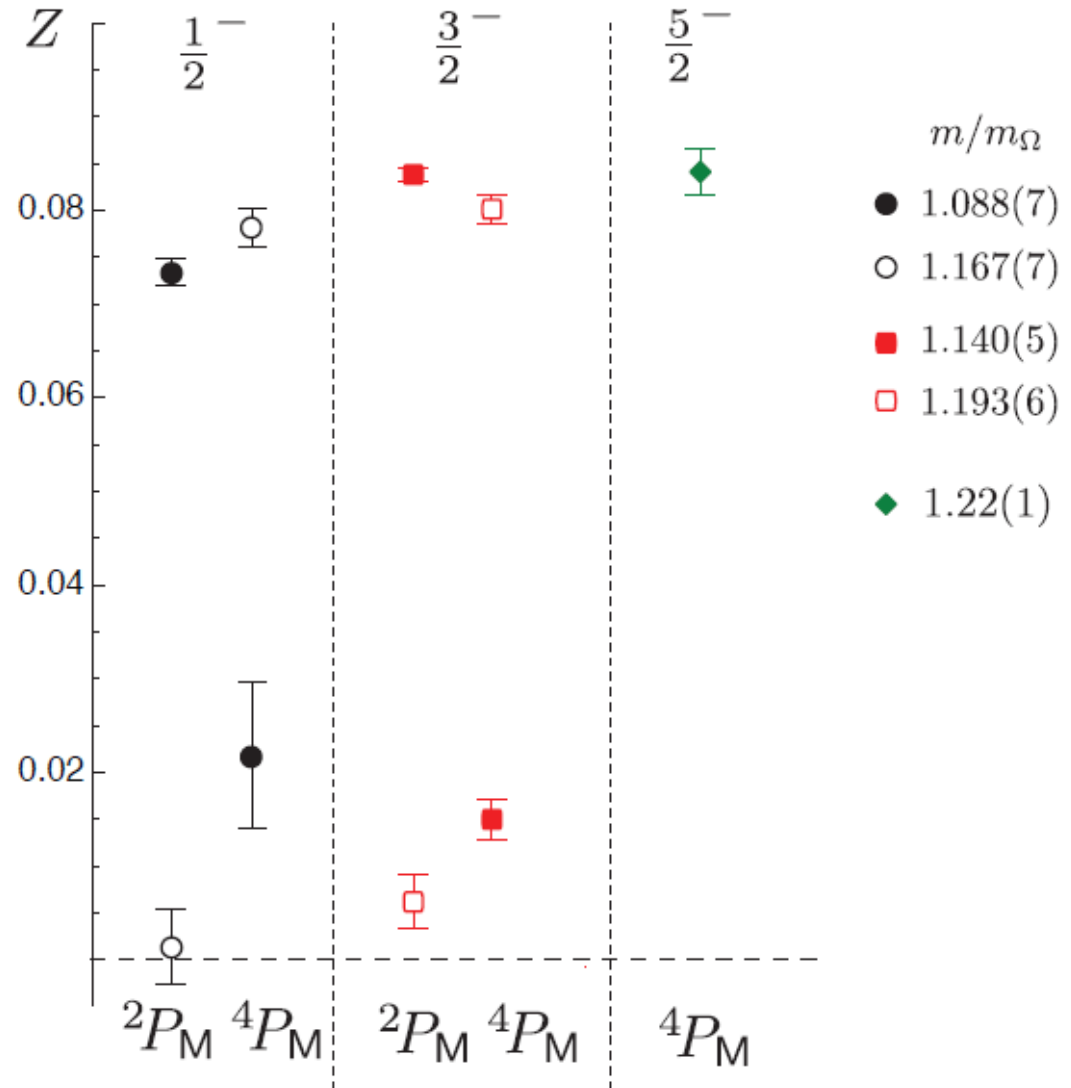


# Nucleon $J^-$

Overlaps

$$Z_i^n = \langle J^- | \Phi_i | 0 \rangle$$

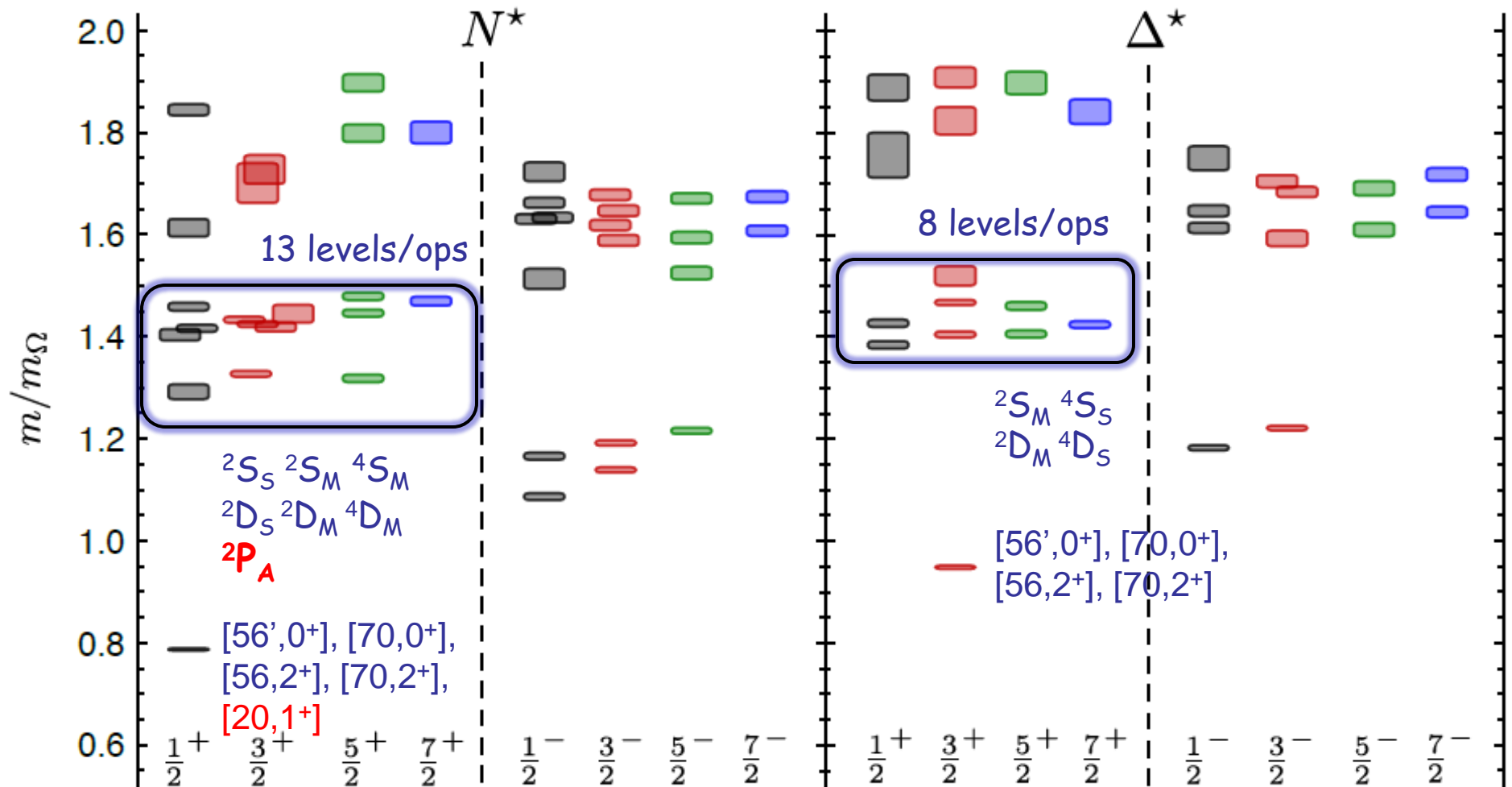
Little mixing in each  $J^-$   
 Nearly "pure" [ $S= 1/2$  &  $3/2$ ]  $1^-$



# N=2 J<sup>+</sup> Nucleon & Delta spectrum

Discern structure: spectral overlaps

Significant mixing in J<sup>+</sup>

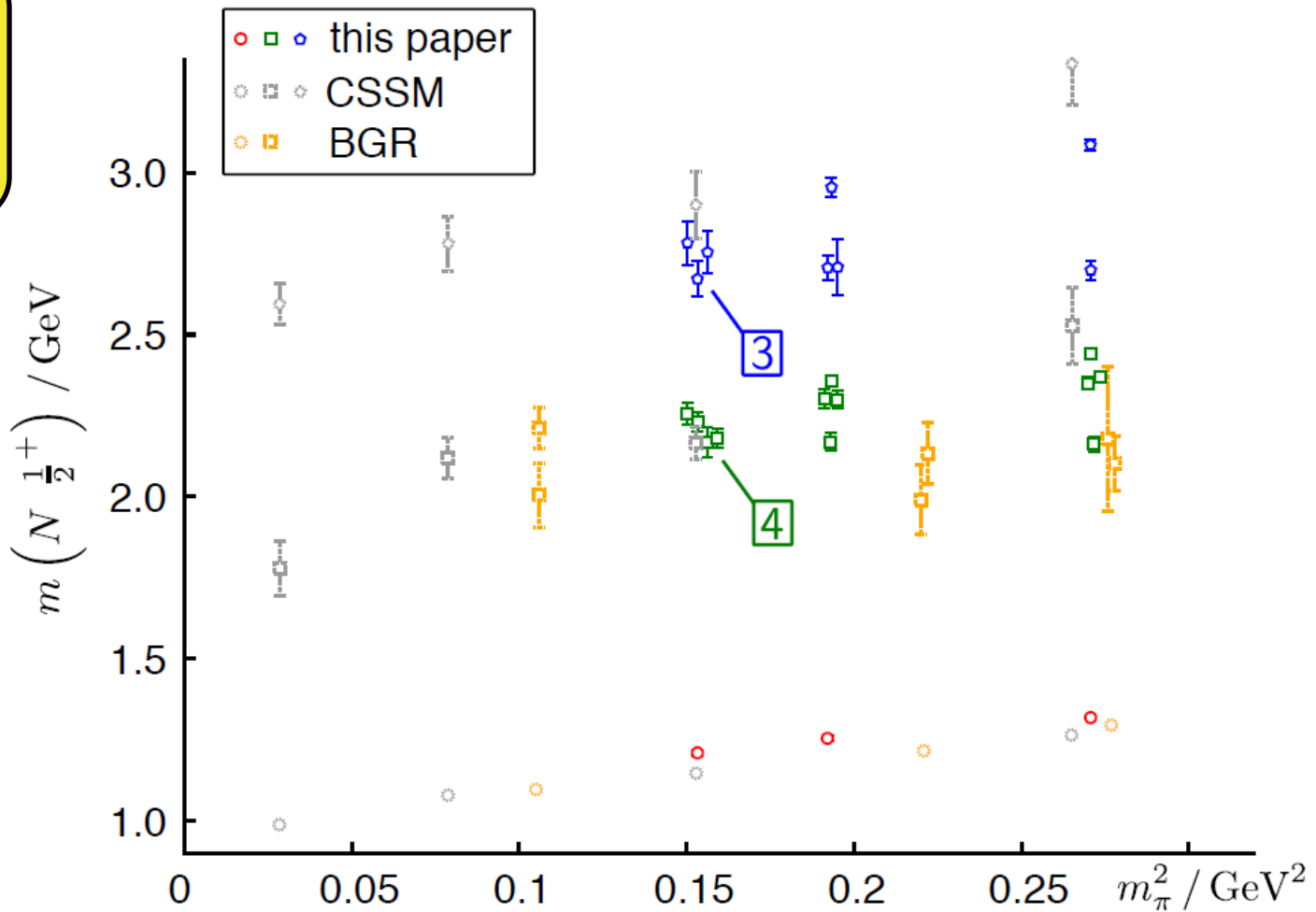


# Roper??

Near degeneracy in  $\frac{1}{2}^+$  consistent with SU(6) O(3) but heavily mixed

Discrepancies??  
Operator basis –  
spatial structure

What else?  
Multi-particle  
operators



# Prospects

- Strong effort in excited state spectroscopy
  - New operator & correlator constructions → high lying states
- Results for baryon excited state spectrum:
  - No “freezing” of degrees of freedom nor parity doubling
  - Broadly consistent with non-relativistic quark model
  - Add multi-particles → baryon spectrum becomes denser
- Short-term plans: **resonance determination!**
  - Lighter quark masses
  - Extract couplings in multi-channel systems

# Backup slides

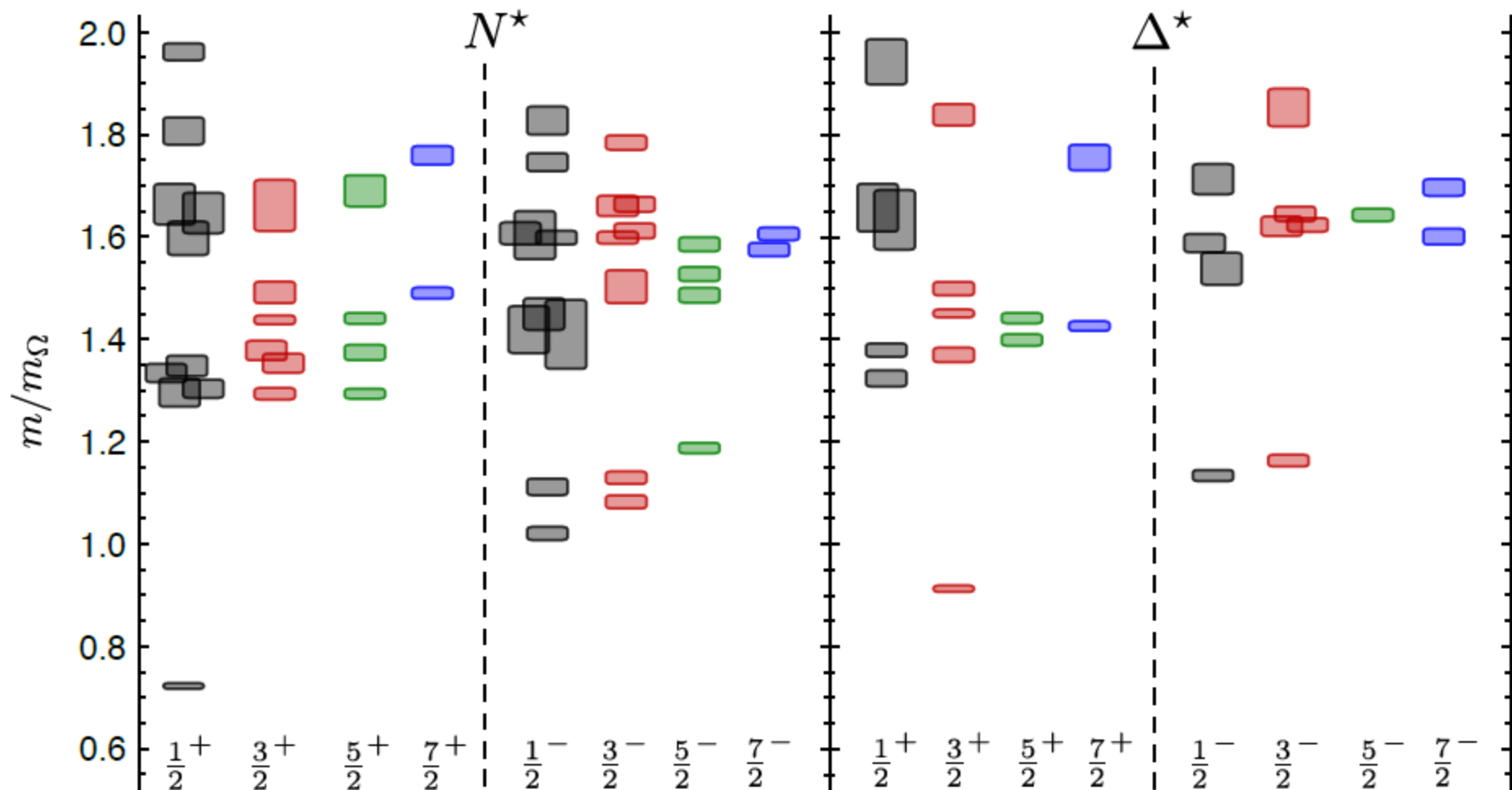
- The end



# N & $\Delta$ spectrum: lower pion mass

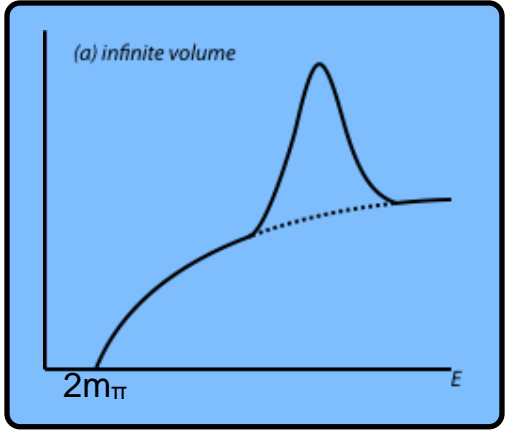
Still bands of states with same counting  
More mixing in nucleon N=2  $J^+$

$m_\pi \sim 400$  MeV

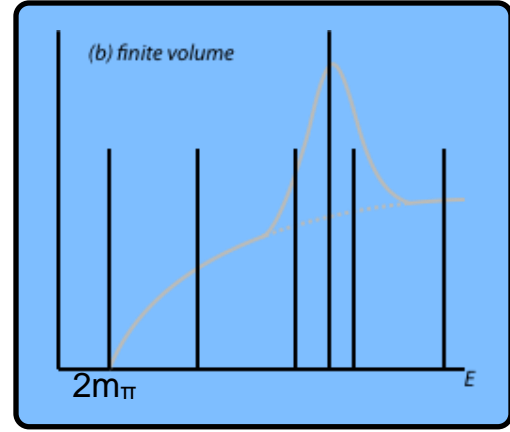


# Spectrum of finite volume field theory

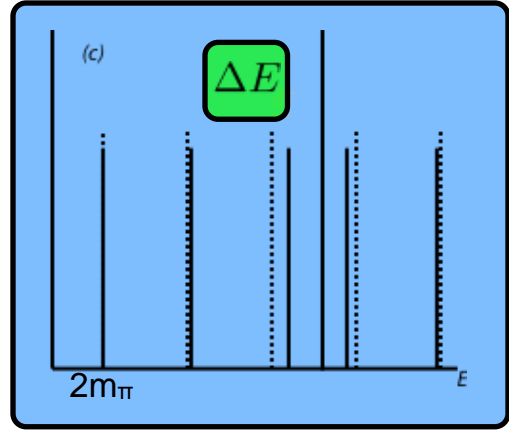
**Missing states:** “continuum” of multi-particle scattering states



**Infinite volume:**  
continuous spectrum  
 $E(p) = 2\sqrt{m_\pi^2 + p^2}$



**Finite volume:** discrete spectrum



Deviation from (discrete) free energies depends upon interaction - contains information about scattering phase shift

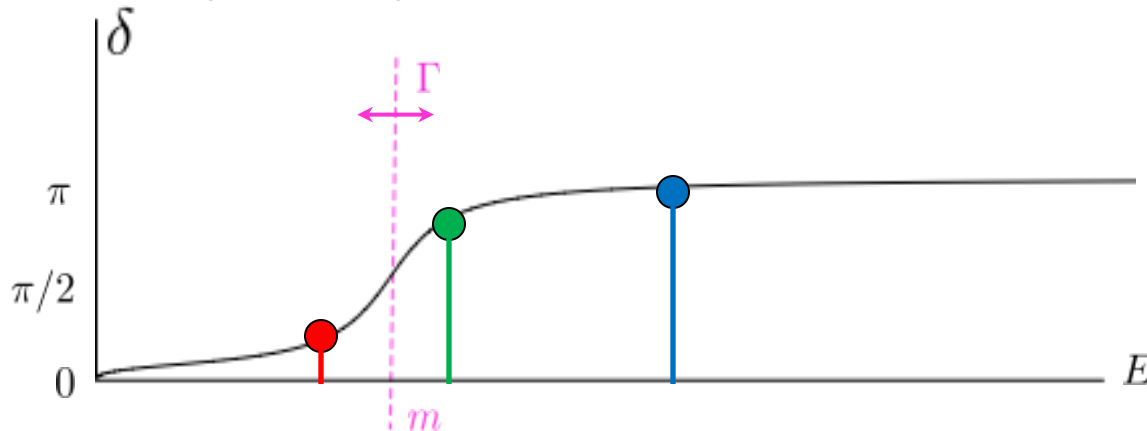
$\Delta E(L) \leftrightarrow \delta(E)$  : Lüscher method

# Finite volume scattering

## Lüscher method

- scattering in a periodic cubic box (length  $L$ )
- finite volume energy levels  $E(L) \rightarrow \delta(E)$

E.g. just a single elastic resonance



e.g.

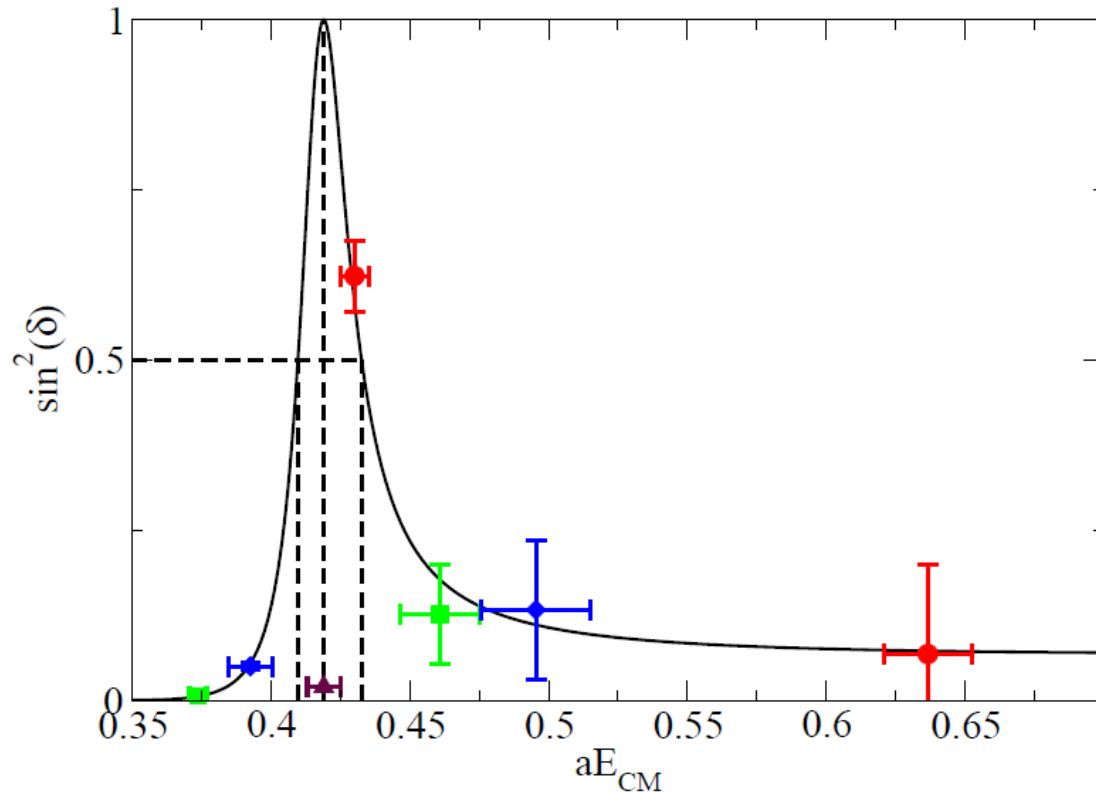
$$\pi\pi \rightarrow \rho \rightarrow \pi\pi$$

$$\pi N \rightarrow \Delta \rightarrow \pi N$$

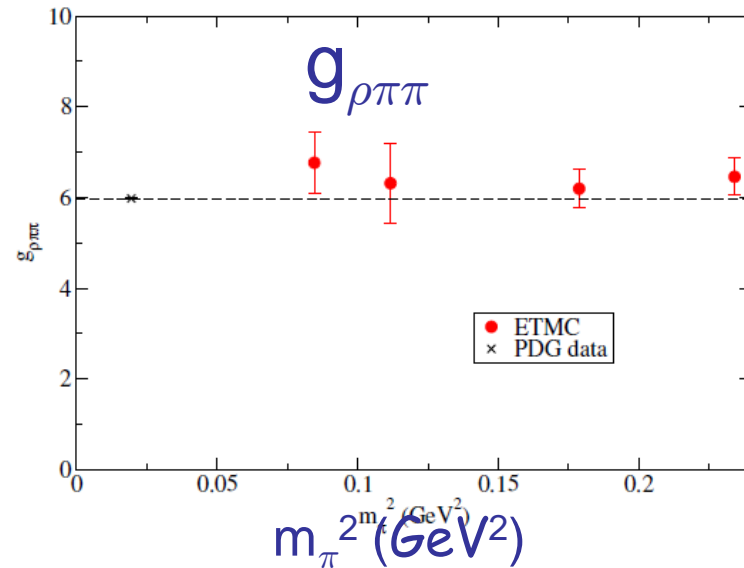
At some  $L$ , have discrete excited energies

# I=1 $\pi\pi$ : the " $\rho$ "

Extract  $\delta_1(E)$  at discrete E



Extracted coupling:  
stable in pion mass



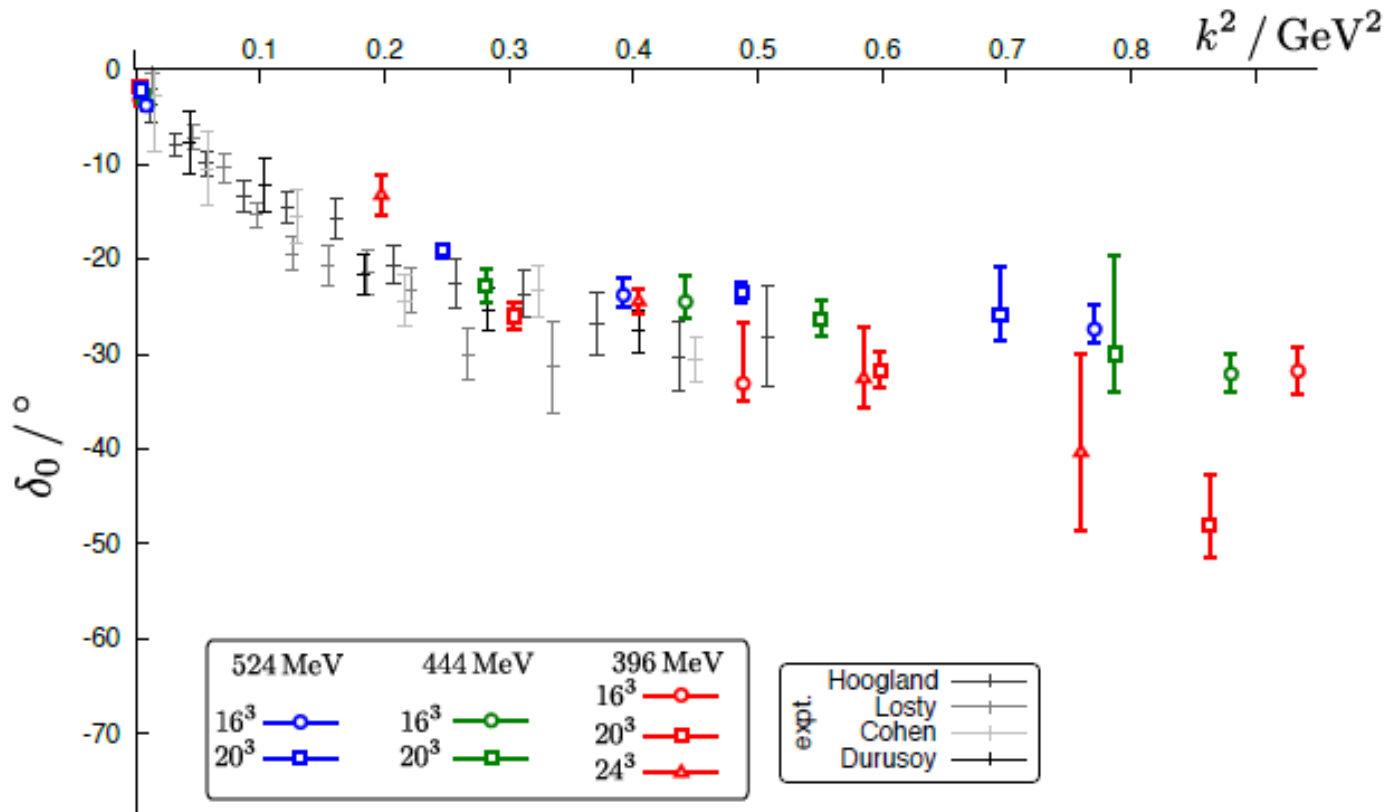
Stability a generic feature  
of couplings??

Feng, Jansen, Renner, 1011.5288

# Phase Shifts demonstration: $I=2 \pi\pi$

$\pi\pi$  isospin=2

Extract  $\delta_0(E)$  at discrete E



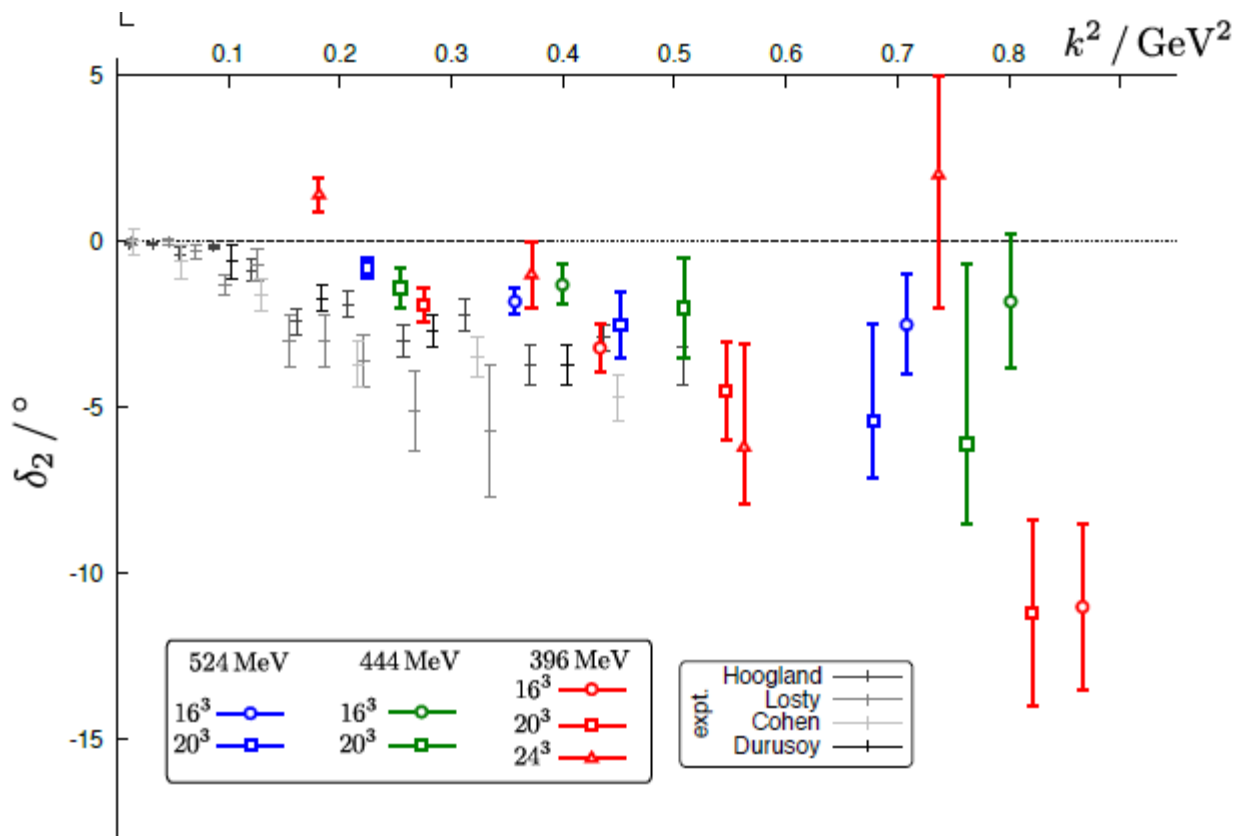
No discernible pion mass dependence

1011.6352 (PRD)

# Phase Shifts: demonstration

$\pi\pi$  isospin=2

$\delta_2(E)$



# Form Factors

What is a form-factor off of a resonance?

What is a resonance? Spectrum first!

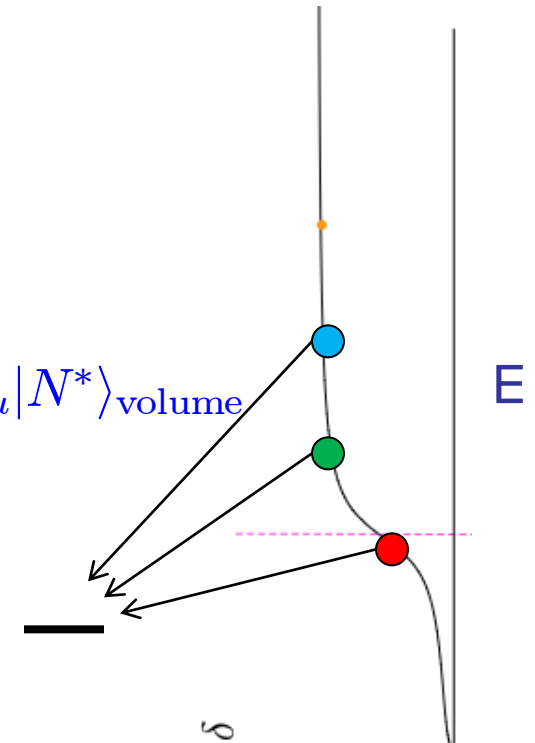
Extension of scattering techniques:

- Finite volume matrix element modified

$$\langle N | J_\mu | N^* \rangle_\infty(Q^2, E) \leftarrow [\delta'(E) + \Phi'(E)] \langle N | J_\mu | N^* \rangle_{\text{volume}}$$

Phase shift

Kinematic factor



Requires excited level transition FF's: some experience

- Charmonium E&M transition FF's (1004.4930)
- Nucleon 1<sup>st</sup> attempt: "Roper" -> N (0803.3020)

Range: few  $GeV^2$

Limitation: spatial lattice spacing

# (Very) Large $Q^2$

Standard requirements:

$$\frac{1}{L} \ll m_\pi, m_N, Q \ll \frac{1}{a}$$

Cutoff effects: lattice spacing  $(a_s)^{-1} \sim 1.6 \text{ GeV}$

Appeal to renormalization group: *Finite-Size* scaling

Use short-distance quantity: compute perturbatively and/or parameterize

$$R(Q^2) = \frac{F(s^2 Q^2)}{F(Q^2)}, \quad s = 2$$

“Unfold” ratio only at low  $Q^2 / s^{2N}$

$$F(Q^2) = R(Q^2/s^2)R(Q^2/s^4) \cdots R(Q^2/s^{2N}) F(Q^2/s^{2N})$$

For  $Q^2 = 100 \text{ GeV}^2$  and  $N=3$ ,  $Q^2 / s^{2N} \sim 1.5 \text{ GeV}^2$

Initial applications: factorization in pion-FF

D. Renner



# Hadronic Decays

Some candidates: determine phase shift  
Somewhat elastic

$m_\pi \sim 400$  MeV

